Modeling erosion in a centrifugal pump in an Eulerian-Lagrangian frame using OpenFOAM®

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Abstract: Erosion induced by solid particle impingement is a very common wear mechanism in turbomachinery and Computational Fluid Dynamics is one of the most widely used tools for its prediction. In this article, erosion is modeled in one of the channels of a centrifugal pump using OpenFOAM®, which is an Open Source CFD package. A review of some of the most commonly used erosion models is carried out in an Eulerian-Lagrangian frame along with a comparative study of the erosion rates obtained with each model. Results yielded some disparities between models due to the different factors taken into consideration. The mesh is then deformed to obtain the resulting eroded geometry.

Keywords: erosion; solid particle impingement; centrifugal pumps

1 Introduction

Computational Fluid Dynamics is one of the most widely used tools for erosion prediction. The use of this method has been enhanced by the major advances in computer technology in recent decades. These allow the attainment of more accurate solutions as well as decreasing computation times significantly. In this work, OpenFOAM’s erosion modeling capabilities are assessed. OpenFOAM is an Open Source CFD package that allows modeling flow conditions in turbomachinery and calculating erosion rates induced by particle impingement in any geometrical configuration. In this article, a single channel of a centrifugal pump is subjected to solid particle impingement and the location of erosion is analyzed and results yielded by four different methods are compared.

2 CFD model

2.1 Particulate phase treatment

Erosion induced by solid particles impinging on complex surfaces is a very common process in many industrial processes. Two different approaches may be taken into account when obtaining numerical solutions to the equations governing fluid flows. The Eulerian-Eulerian method is suitable for high particle concentrations. In this case the particles are treated as a continuous phase and the eulerian continuum equations are solved for the two different phases. When the concentration of the particles inside the fluid is sufficiently low, the continuous phase would be solved in the first place, while the successive positions and velocities of the particles would be integrated from Newton’s equations for motion. This would be the so called Eulerian-Lagrangian approach. The influence of the particulate phase on the continuous phase can be neglected when the concentration of the disperse phase is adequately low. Calculation of the Particle Mass Loading [1] is a convenient tool for making the correct choice of methodology.

In the case being presented in this article, no significant two-way particle-fluid coupling is to be expected due to the calculated particle mass loading being slightly over 0.015. This indicates that the momentum transfer from the particles to the continuous phase is negligible, as it can be verified in the simulations by inspecting the recorded values of this momentum transfer.
Figure 1: Back-side of the centrifugal pump impeller volume.

Figure 2: Centrifugal pump channel discretization and boundaries.

2.2 Domain discretisation

For the present article, a random example of a centrifugal pump impeller has been chosen, and its corresponding volume is presented in Figure 1.

In order to keep the computational cost to a minimum while maintaining a reasonable accuracy, only one channel of the impeller has been taken into consideration, along with the proportional part of the inlet boundary. Figure 2 shows the geometrical configuration of the channel, which has been divided into several parts for an optimized discretization procedure and the location of the different boundary conditions. In the case of the inlet boundary condition, a uniform velocity was chosen, no slip was chosen for the walls of the impeller and pressure outlet for the outlet boundary. As it can be observed in Figure 2, the throat bush was sliced in order to include its proportional part in the simulation. It is for that reason that cyclic boundary conditions were chosen at the sides of the inlet.

In this case, the mesh and procedure for the calculation of erosion have been implemented so that the time needed to obtain the contours of erosion and the resulting deformed geometry is minimized. The mesh chosen is composed almost entirely by hexahedral cells. Given that the turbulence model to be employed is the k-omega, the \( y^+ \) was also calculated. The values obtained indicate the suitability of the mesh for this case, ranging from 1 to 60, which was considered adequate for the model used and for the calculation of the particle variables. Quality parameters for the mesh were ensured by the fact that the mesh is composed almost in its entirety by hexahedral cells.

3 Steady state solution and Implementation of a Simple Reference Frame solver for an incompressible phase and Lagrangian particles

3.1 Steady state results

The first step in the present methodology is the attainment of the steady state for the fluid phase. This was achieved by selecting the appropriate solver from the available ones in OpenFoam: SRFSimpleFoam. This solver is suitable for rotating domains on a Single Reference Frame (SRF). The discretization schemes used were second order for the velocity and first order for the k-omega turbulence modeling. Final solution for the steady state was achieved when residuals fell below \( 1 \times 10^{-4} \). The resulting velocity field is represented in Figure 3, which shows a slice of the channel at its central part. The rotational velocity of the pump was 2850 revolutions per minute and the flow rate at the inlet was 0.0063 m\(^3\)/h.

Figure 3: Half of the channel showing steady state velocity contours.
3.2 Transient simulation of the particulate phase

OpenFOAM's customisability is one of its main incentives, together with the ease of manipulation of the different libraries and the linkage possibilities between them. For the case being treated here, the lagrangian-intermediate library is chosen for coupling with the incompressible solver. This library contains a series of templated functions, models and forces that allow calculation of the particle variables once the carrier phase has been previously solved. The intermediate library provides the user with five different classes for representing the particles. As no particle-particle collisions, chemical reactions or temperature effects are taken into consideration; the Kinematic class is the most suitable library to be coupled to the incompressible solver. The Kinematic library allows two different particle treatments. The first one is to track the particles individually, while the second one avoids extra computational costs by tracking a set of particles, or is also called computational parcel. In a computational parcel, in order to capture the behaviour of the real particles, some real case properties are defined. For the case being taken into consideration, the individual particle tracking approach is selected, due to the previously mentioned low concentration of particles in the fluid phase.

The selection of the incompressible solver was SRF-PimpleFoam, which is a transient solver for incompressible flow in a single rotating frame.

The procedure starts with the solution of the fluid flow. Once the continuous phase solution is calculated, the different forces on the particles are computed. Newton's equations for motion are then solved and the position of the particles is obtained by means of integration. The difference between the Simple Reference Solver and the pimpleFoam incompressible solver lies in the treatment of the velocity. The fluid's velocity inside the rotating volume is calculated as relative, and the particles are affected by an SRF Force (due to the rotating component of the movement) which has to be incorporated into the disperse phase properties. An additional momentum source to take into account the drag force exerted by the fluid is included.

Linking the lagrangian library to the incompressible solver means adding the constructors and files needed to represent the particles inside the fluid. Then inclusion inside the solver of the required functions that will take the continuous phase solution's results in order to apply them to the particles is performed. These functions are responsible for the calculation of the forces on the particles before the mentioned integration of the acceleration in order to obtain their positions and velocities.

4 Erosion modeling

4.1 Erosion calculation

The Kinematic library is largely based on templates, which is a powerful feature of C++ programming language. This language requires declaring variables, functions and most kinds of entities using types. Templates are functions or classes that are defined for one or more types that are not yet specified. Due to the fact that a lot of parts of the code look the same for different types, templates become a very useful tool and that is the reason why they are very common within OpenFOAM's code. The function that calculates erosion is in fact, a template, which in this case is defined for the Kinematic class. Impingement information, such as impact speed and impact angle, is gathered as the individual particles or computational parcels hit the specified walls of the geometry on which erosion needs to be calculated. After this, the recorded information is used to calculate erosion at each particular face of the boundary and the field formed by all the faces is stored to a file for post processing. The erosion field is calculated at each face of the boundaries by means of the addition of each particle's calculated induced erosion. At the moment each particle reaches the specified boundaries, the face that was hit is recorded and erosion is calculated and added together with the previous contributions and stored in the erosion field.

4.2 Implemented erosion models

Erosion induced by solid particles carried by a fluid is a common deterioration mechanism that takes place in many different environments and kinds of equipment. It is because of this that the literature that can be found on the subject is vast. However, the level of agreement between the different authors is very low. Meng and Ludema [2] carried out a literature review of more than 5000 articles published in the most important erosion related journals and conference proceedings. Within these, they encountered almost 2000 erosion models, from which they separated 28 different ones. Discrepancies between different authors may be explained by pointing out the extreme complexity of erosion processes and the number of variables involved in its calculation. This complexity was illustrated by J.A.C.
Humphrey in [3], where he listed all the factors affecting erosion that had been previously investigated and grouped them in three different categories, namely, the ones affecting the particles, the surfaces involved in the process and the fluid flow. In fact, one of the only theories on which there seems to be some kind of agreement is the one differentiating between two different mechanisms affecting the process: Cutting and Deformation Wear. Definition of these mechanisms has been adopted by numerous authors ([4–9]).

This theory states that, on one side, particles may hit the surface and tear material away with them in a cutting manner. This is called cutting wear and it is the predominant mechanism for ductile materials and particles impinging at low angles. On the other side, several particles might impact on the same place transferring some of their kinetic energy to the surface in the form of hardening work. Eventually, a piece of material detaches from the surface being subjected to erosion and is carried away by the fluid. This mechanism is called deformation wear and it prevails for high angles of impingement and brittle materials.

In this article, in addition to the built-in erosion model [10], two additional equations have been implemented in OpenFOAM. Given that each equation is the result of a personal approach (in fact, literally thousands of different models can be encountered in the literature) only the location of erosion will be compared, paying special attention to the location of the maximum erosion. Calculation of wear rates for a particular case would imply experimental tests to obtain some of the constants used in the models. However, if the location of maximum erosion and the relative difference between the different eroded parts is successfully represented by the model, the wear experienced by the geometry can be obtained by scaling that obtained values to fit the available results.

It is worth noting that the four different models calculate the wear rates for exactly the same number of impact locations, angles and velocities. This is due to the fact that the information concerning each impact is simultaneously given as input to the four erosion models. Due to the very different nature of the equations used for erosion calculations, the most significant data we can take from the calculations is the location of the maximum erosion along with the relative difference in erosion between the different areas of the channel. It is for this reason that the labeling only specifies maximum and minimum erosion locations. Once we have the real geometry to compare our results with, if the locations and the relative difference match up, it is only a matter of scaling the numerical values to fit the real ones.

### 4.2.1 Finnie Erosion model

The present model is one of the first and still widely used models in literature [10]. It is divided into two different equations, which depend on the angle of impingement. These equations yield the volume of material removed by a single abrasive particle of mass \( m \), travelling at a velocity denoted by \( V \). Two experimental constants are present in the set of equations, namely, the ratio of depth of contact to the depth of cut and the ratio of vertical to horizontal force components. The material’s plastic flow stress of the material subjected to erosion and the angle of impingement are the last two variables to complete the formulae.

It is worth underlining that this model yields accurate results for low angles of impingement (which is where maximum erosion usually takes place in ductile metals) but it predicts no erosion at normal incidence. In a second publication [11], Finnie modified his form in order to account for some other variables that would allow capturing some erosion for normal incidence. Calculated results for Finnie’s model are shown in Figure 4 and the model is shown in Equations (1) and (2).

The equation yields the volume of material, \( Q \) removed by a single abrasive grain of mass, \( m \), and velocity \( V \):

\[
Q = \frac{mV^2}{p\psi K} \left( \sin 2\alpha - \frac{6}{K} \sin^2 \alpha \right) \quad \text{if} \quad \tan \alpha \leq \frac{K}{6} \tag{1}
\]

\[
Q = \frac{mV^2}{p\psi K} \left( \frac{K \cos^2 \alpha}{6} \right) + \quad \text{if} \quad \tan \alpha > \frac{K}{6},
\]

where \( p \) is the plastic flow stress of the material, \( \alpha \) is the angle of impingement, \( \psi \) is a constant that represents the ratio of depth of contact to the depth of cut and \( K \) is the ratio of vertical to horizontal force components on the particle.

### 4.2.2 Grant and Tabakoff Erosion model

The main assumption made for the development of this equation is that erosion depends on two different mechanisms that take place at low and high angles of impingement respectively. A combination of these two mechanisms is used when calculating erosion for intermediate angles of impingement.

Obtained erosion with Grant and Tabakoff’s model [9] can be seen in Figure 5 and the implemented model is...
Figure 4: Finnie’s model calculated erosion.

Figure 5: Grant and Tabakoff’s model calculated erosion.

Figure 6: Nandakumar’s model calculated erosion.

shown in Equation (3).

\[ \varepsilon = K_1 f(\beta_1) \left( V_{IT}^3 - V_{IT}^2 \right) + f(V_{1N}), \]  

(3)

where

\( \varepsilon \) = the erosion damage per unit mass of impacting particles

\( K_1 \) = material contact

\( f(\beta_1) \) = empirical function of particle impact angle

\( V_{IT} \) = tangential component of incoming particle velocity

\( V_{IT} \) = tangential component of rebounding particle velocity

\( f(V_{1N}) \) = component of erosion due to the normal component of velocity

4.2.3 Nandakumar erosion model

The last model was developed by Nandakumar et al. [12] who took as a starting point other authors work (Finnie [10] and Bitter [5, 6]). This model calculated the volume eroded by the impacts of the particles on the material’s surface. Again empirical constants are present in this model so that experimental data is needed in order to calculate the eroded volume for a particular case.

Figure 6 shows the location and relative difference of the erosion predicted by Nandakumar’s model and it’s mathematical expression is detailed in Equation (4).

\[ \Delta Q = C m p_0^{0.15} (V_0 \sin \Theta)^{2.3} \]

\[ D m^{1.1875} d_p^{-0.0625} V_0^{2.375} (\cos \Theta)^3 (\sin \Theta)^{0.375}, (4) \]

where \( C \) and \( D \) are empirical constants, \( m \) is the mass of the particle, \( p \) is the density of the particle, \( \Theta \) is the angle of impact, \( V_0 \) is the impact velocity and \( d_p \) is the diameter of the particle.

5 Results discussion

It is clear that the three different models predict the location of the maximum erosion at the same location. However, discrepancies arise when quantifying that erosion and the relative erosion around that maximum erosion location. While Finnie’s and Grant and Tabakoff’s’s erosion models both yield similar results for the considered pump channel, Nandakumar’s model predicts wear in locations where the other two don’t. Given that the three models yield similar results for the location of the maximum erosion, it seems reasonable to choose Hashish’s model for a first approach, due to the lack of experimental constants in the formulæ being thus, not necessary to perform experimental tests to obtain them.

However for a time dependent evolution of the wear scar, a much more complex approach would have to be considered, applying surface deformation so that the evolution of the velocities and angles of impingement with the development of the eroded surface are properly accounted for.
Modeling erosion in a centrifugal pump

The work to be continued in the future includes the study of the flow changes with surface deformation due to erosion, thanks to the development of the deformation algorithm commented here along with the validation of the CFD model for low concentrations by means of Particle Image Velocimetry as well as the implementation of a methodology for erosion prediction that accounts for the surface changes with time.

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References


Figure 7: Pump channel close up before deformation.

Figure 8: Deformed pump channel according to erosion.

6 Geometry deformation example

Figure 7 shows the initial shape of the pump channel while Figure 8 shows the same channel after the deformation algorithm that the authors developed in OpenFoam is applied. Finally, Figure 8 provides confirmation that the deformation of the surface follows the erosion contours, i.e., the higher the erosion rate, the larger the deformation. Deformation of the surface is the first step towards a more accurate calculation of how erosion develops, which is one of the next steps that the authors intend to follow.

7 Conclusions and future work

Provided that an enormous number of different erosion models can be found in the literature, it seems obvious that one could search for a particular method suitable for the application taken into consideration. However, this search might be time consuming. If only a reasonable initial approach is needed, Hashish’s model appears to be the best choice, while for more accurate predictions a literature review should be performed searching for a more suitable method for each particular application.